

Symposium on Geometry Processing – Paris – 2018 Course on Numerical Optimal Transport – Bruno Lévy

OVERVIEW

Part. 1. Goals and Motivations

Part. 2. Introduction to Optimal Transport

Part. 3. Semi-Discrete Optimal Transport

Part. 4. Applications in Computational Physics



Goals and Motivations

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Goal #1: "Understanding"



Part. 1 Optimal Transport Goal #1: "Understanding" What I can't create I don't understand **Richard Feynman**



Part. 1 Optimal Transport Goal #1: "Understanding"





Jos Stam, Stable Fluids, 1999 The art of fluid sim.

Understand fluids Explain Program



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Goal #1: "Understanding"

+1129+5

I have no formal background in fluid dynamics. I am not an engineer nor do I have a specialized degree in the mathematics or physics of fluids. I am fortunate that I did not have to carry that baggage around. On the other hand, I *do* have degrees in pure mathematics and computer science and have an artsy background. More importantly, I have written computer code that animates fluids.*

I wrote code That is the bottom line.

I wrote code



Goal #1: "Understanding" Your mission statement:

- 1. Understand the stuff
- 2. Explain it in simple terms

Be a good teacher, to others and to yourself Know what you know and what you don't know Try to know what you don't know

3. Program it



Part. 1 Optimal Transport Measuring distances between functions





Part. 1 Optimal Transport Measuring distances between functions





Part. 1 Optimal Transport Measuring distances between functions





$$d_{L_2}(f_1, f_2) = \int \left(f_1(x) - f_2(x) \right)^2 dx$$



Part. 1 Optimal Transport Measuring distances between function





$$d_{L_2}(f_1, f_2) = \int \left(f_1(x) - f_2(x) \right)^2 dx$$





















































Part. 1 Optimal Transport Gaspard Monge - 1784

666. MÉMOIRES DE L'ACADÉMIE ROYALE

MÉMOIRE SURLA THÉORIE DES DÉBLAIS ET DES REMBLAIS.

Par M. MONGE.

L'orsqu'on doit transporter des terres d'un lieu dans un autre, on a coutume de donner le nom de *Déblai* au volume des terres que l'on doit transporter, & le nom de







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Part. 1 Optimal Transport Gaspard Monge – geometry and light











Part. 1 Optimal Transport Monge-Brenier-Villani, the french connection



Cédric Villani Optimal Transport Old & New Topics on Optimal Transport



Yann Brenier The polar factorization theorem (Brenier Transport)

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Optimal transport geometry and light

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[Caffarelli, Kochengin, and Oliker 1999]







[Castro, Merigot, Thibert 2014]

Part. 1 Optimal Transport – Image Processing





Barycenters / mixing textures

[Nicolas Bonneel, Julien Rabin, Gabriel Peyre, Hanspeter Pfister] Video-style transfer, A.I., "data sciences"

[Nicolas Bonneel, Kalyan Sunkavalli, Sylvain Paris, Hanspeter Pfister] [Marco Cuturi, Gabriel Peyré]



Optimal transport - geometry and light



[Chwartzburg, Testuz, Tagliasacchi, Pauly, SIGGRAPH 2014]





Part. 1. Motivations

Discretization of functionals involving the Monge-Ampère operator, Benamou, Carlier, Mérigot, Oudet arXiv:1408.4536

The variational formulation of the Fokker-Planck equation Jordan, Kinderlehrer and Otto SIAM J. on Mathematical Analysis









How to "morph" a shape into another one of same mass while minimizing the "effort" ?





How to "morph" a shape into another one of same mass while minimizing the "effort" ?

The "effort" of the best T defines a **distance** between the shapes





How to "morph" a shape into another one while preserving mass and minimizing the effort ?







How to "morph" a shape into another one while preserving mass and **minimizing the effort**?

"minimum action principle"




How to "morph" a shape into another one while **preserving mass** and **minimizing the effort** ? "conservation law" "minimum action principle"



OT=

"minimum action principle subject to conservation law"

Yann Brenier: "Each time the Laplace operator is used in a PDE, it can be replaced with the Monge-Ampère operator"

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OT=

"minimum action principle subject to conservation law"

Yann Brenier: "Each time the Laplace operator is used in a PDE, it can be replaced with the Monge-Ampère operator"

New ways of simulating physics with a computer



OT=

"minimum action principle subject to conservation law"

Yann Brenier: "Each time the <u>Laplace operator</u> is used in a PDE, it can be replaced with the Monge-Ampère operator"

<u>Fast Fourier Transform</u>

New ways of simulating physics with a computer

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OT=

"minimum action principle subject to conservation law"

Yann Brenier: "Each time the Laplace operator is used in a PDE, it can be replaced with the Monge-Ampère operator" Fast Fourier Transform

New ways of simulating physics with a computer

Fast OT algo. ???

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Optimal Transport an elementary introduction







(X;µ)

(Y;v)

Two measures
$$\mu$$
, v such that $\int_X d\mu(x) = \int_Y dv(x)$







(X;µ)

(Y;v)

A map T is a *transport map* between μ and ν if $\mu(T^{-1}(B)) = \nu(B)$ for any Borel subset B of Y

(Borel subset = subset that can be measured)







(X;µ)



A map T is a *transport map* between μ and \vee if $\mu(T^{-1}(B)) = \nu(B)$ for any Borel subset B of Y









A map T is a *transport map* between μ and \vee if $\mu(T^1(B)) = \vee(B)$ for any Borel subset B of Y







(X;µ)

(Y;v)

A map T is a *transport map* between μ and ν if $\mu(T^1(B)) = \nu(B)$ for any Borel subset B of Y

Notation: if T is a *transport map* between μ and ν then one writes $\nu = T#\mu$ (ν is the *pushforward* of μ)











Monge's problem:

Find a transport map T that minimizes $C(T) = \int_X ||x - T(x)||^2 d\mu(x)$



Monge's problem:

Find a transport map T that minimizes $C(T) = \int_X ||x - T(x)||^2 d\mu(x)$

- Difficult to study
- If μ has an atom (isolated Dirac), it can only be mapped to another Dirac (T needs to be a map)



Monge's problem:

Find a transport map T that minimizes $C(T) = \int_X ||x - T(x)||^2 d\mu(x)$



Transport from a measure concentrated on L₁ onto another one concentrated on L₂ and L₃



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Transport from a measure concentrated on L₁ onto another one concentrated on L₂ and L₃

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Monge's problem:

Find a transport map T that minimizes $C(T) = \int_X ||x - T(x)||^2 d\mu(x)$



Transport from a measure concentrated on L₁ onto another one concentrated on L₂ and L₃

The infimum is never realized by a map, need for a relaxation

Monge's problem:

Find a transport map T that minimizes $C(T) = \int_X ||x - T(x)||^2 d\mu(x)$

Kantorovich's problem (1942):

Find a measure γ defined on X x Y such that $\int_{X \text{ in } X} d\gamma(x,y) = d\nu(y)$ and $\int_{Y \text{ in } Y} d\gamma(x,y) = d\mu(x)$

that minimizes
$$\iint_{X \times Y} || x - y ||^2 d_{Y(x,y)}$$





Monge's problem:

Find a transport map T that minimizes $C(T) = \int_X ||x - T(x)||^2 d\mu(x)$

Kantorovich's problem:

Find a measure γ defined on X x Y such that $\int_{X \text{ in } X} d\gamma(x,y) = d\nu(y)$ and $\int_{Y \text{ in } Y} d\gamma(x,y) = d\mu(x)$

 $``\gamma(x,y)"$: How much sand goes from x to y

that minimizes $\iint_{X \times Y} || x - y ||^2 d_{\gamma(x,y)}$



Monge's problem:

Find a transport map T that minimizes $C(T) = \int_X ||x - T(x)||^2 d\mu(x)$

Kantorovich's problem:

Find a measure
$$\gamma$$
 defined on X x Y
such that $\int_{X \text{ in } X} d\gamma(x,y) = d\nu(y)$
and $\int_{Y \text{ in } Y} d\gamma(x,y) = d\mu(x)$
that minimizes $\iint_{X \times Y} ||x - y||^2 d\gamma(x,y)$

Monge's problem:

Find a transport map T that minimizes $C(T) = \int_X ||x - T(x)||^2 d\mu(x)$

Kantorovich's problem:

Find a measure
$$\gamma$$
 defined on X x Y
such that $\int_{X \text{ in } X} d\gamma(x,y) = dv(y)$
and $\int_{Y \text{ in } Y} d\gamma(x,y) = d\mu(x)$
Everything that is
transported to y sums to "v(y)"

that minimizes $\iint_{X \times Y} || x - y ||^2 d_{Y(x,y)}$





Transport plan – example 1/4 : translation of a segment





Transport plan – example 1/4 : translation of a segment





Transport plan – example 2/4 : spitting a segment



Observation 1. If $(Id \times T) \sharp \mu \in \pi(\mu, \nu)$, then T pushes μ to ν .



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Proof. $(Id \times T) \sharp \mu$ belongs to $\pi(\mu, \nu)$, therefore $(P_2) \sharp (Id \times T) \sharp \mu = \nu$, or $((P_2) \circ (Id \times T)) \sharp \mu = \nu$, thus $T \sharp \mu = \nu$ \Box



Observation 1. If $(Id \times T) \sharp \mu \in \pi(\mu, \nu)$, then T pushes μ to ν .

Proof. $(Id \times T) \sharp \mu$ belongs to $\pi(\mu, \nu)$, therefore $(P_2) \sharp (Id \times T) \sharp \mu = \nu$, or $((P_2) \circ (Id \times T)) \sharp \mu = \nu$, thus $T \sharp \mu = \nu$ \Box

With this observation, for transport plans of the form $\gamma = (Id \times T) \sharp \mu$, (K) becomes

$$\min\left\{\int_{\Omega\times\Omega} c(x,y)d\left((Id\times T)\sharp\mu\right)\right\} = \min\left\{\int_{\Omega} c(x,T(x))d\mu\right\}$$





Transport plan – example 3/4 : splitting a Dirac into two Diracs





Transport plan – example 3/4 : splitting a Dirac into two Diracs

(No transport map)





Transport plan – example 4/4 : splitting a Dirac into two segments





Transport plan – example 4/4 : splitting a Dirac into two segments

(No transport map)





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Duality is easier to understand with a discrete version Then we'll go back to the continuous setting.





(DMK): Min <C, γ > $s.t. \begin{cases} P_1 \gamma = u \\ P_2 \gamma = v \\ \gamma \ge 0 \end{cases}$





(DMK): Min <C, γ > s.t. $\begin{cases}
P_1 \gamma = u \\
P_2 \gamma = v \\
\gamma \ge 0
\end{cases}$







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Part. 2 Optimal Transport – Duality

< *u*, *v* > denotes the dot product between *u* and *v*

(DMK):

Consider $\mathcal{I}(\phi, \psi) = \langle c, \gamma \rangle - \langle \phi, P_1 \gamma - u \rangle - \langle \psi, P_2 \gamma - v \rangle$



Part. 2 Optimal Transport – Duality(DMK):
Min <C, γ >
S.t.Min <C, γ >
P₁ γ = U
P₂ γ = V
 $\gamma \ge 0$

Consider
$$\mathcal{I}(\phi, \psi) = \langle c, \gamma \rangle - \langle \phi, P_1 \gamma - u \rangle - \langle \psi, P_2 \gamma - v \rangle$$

$\begin{array}{l} \mbox{Remark: Sup[} \ \mathcal{I}(\phi,\psi) \ \] = < c, \ \gamma > \ if \ P_1 \ \gamma = u \ and \ P_2 \ \gamma = v \\ \phi \ \in \ \mathrm{IR}^m \\ \psi \in \ \mathrm{IR}^n \end{array}$



Part. 2 Optimal Transport – Duality(DMK):
Min <C, γ >
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Part. 2 Optimal Transport – Duality(DMK):
Min <C, γ >
S.t.Min <C, γ >
P₁ γ = U
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 $\gamma \ge 0$

Consider
$$\mathcal{I}(\phi, \psi) = \langle c, \gamma \rangle - \langle \phi, P_1 \gamma - u \rangle - \langle \psi, P_2 \gamma - v \rangle$$

 $\begin{array}{ll} \mbox{Remark: Sup[} \ \mathcal{I}(\phi,\psi) \] = < c, \ \gamma > \ \mbox{if } P_1 \ \gamma = u \ \mbox{and } P_2 \ \gamma = v \\ & \phi \ \in \ \mbox{IR}^m \\ & \psi \ \in \ \mbox{IR}^n \end{array} = + \infty \ \mbox{otherwise} \end{array}$

Consider now: Inf $\begin{bmatrix} Sup[\mathcal{I}(\phi, \psi) \end{bmatrix} \end{bmatrix}$ $\gamma \ge 0 \quad \begin{array}{c} \phi \in IR^m \\ \psi \in IR^n \end{array}$



Part. 2 Optimal Transport – Duality(DMK):
Min \gamma>
S.t.
$$P_1 \gamma = u$$

 $P_2 \gamma = v$
 $\gamma \ge 0$

Consider
$$\mathcal{I}(\phi, \psi) = \langle c, \gamma \rangle - \langle \phi, P_1 \gamma - u \rangle - \langle \psi, P_2 \gamma - v \rangle$$

 $\begin{array}{ll} \mbox{Remark: Sup[} \ \mathcal{I}(\phi,\psi) \] = < c, \ \gamma > \mbox{if } P_1 \ \gamma = u \ and \ P_2 \ \gamma = v \\ & \phi \ \in \ IR^m \\ & \psi \ \in \ IR^n \end{array} = +\infty \ otherwise \end{array}$



Part. 2 Optimal Transport – Duality(DMK):
Min \gamma>
S.t.
$$\begin{cases} P_1 \ \gamma = u \\ P_2 \ \gamma = v \\ \gamma \ge 0 \end{cases}$$

Consider
$$\mathcal{I}(\phi, \psi) = \langle c, \gamma \rangle - \langle \phi, P_1 \gamma - u \rangle - \langle \psi, P_2 \gamma - v \rangle$$

Remark: Sup[$\mathcal{I}(\phi, \psi)$] = < c, γ > if P₁ γ = u and P₂ γ = v $\phi \in IR^{m}$ $\psi \in IR^{n}$ = + ∞ otherwise

 $\begin{array}{ll} \text{Consider now: Inf} \left[\begin{array}{c} \text{Sup} \left[\begin{array}{c} \mathcal{I}(\phi,\psi) \end{array} \right] \right] = \text{Inf} \left[\begin{array}{c} < c, \ \gamma > \end{array} \right] & (\text{DMK}) \\ \gamma \geq 0 & \phi \in \mathrm{IR}^m & \gamma \geq 0 \\ \psi \in \mathrm{IR}^n & P_1 \ \gamma = u \\ P_2 \ \gamma = v \end{array} \end{array}$



Part. 2 Optimal Transport – Duality(DMK):
Min \gamma >
S.t.
$$\begin{cases} P_1 \ \gamma = u \\ P_2 \ \gamma = v \\ \gamma \ge 0 \end{cases}$$
Inf $\begin{bmatrix} Sup[< C, \gamma > - < \phi, P_1 \ \gamma - u > - < \psi, P_2 \ \gamma - v >] \end{bmatrix}$

Part. 2 Optimal Transport – Duality(DMK):
Min \gamma >
S.t.
$$\begin{cases} P_1 \gamma = u \\ P_2 \gamma = v \\ \gamma \ge 0 \end{cases}$$
Inf $\begin{bmatrix} Sup[- <\phi, P_1 \gamma - u > - <\psi, P_2 \gamma - v >] \end{bmatrix}$ $\gamma \ge 0$ $\psi \in IR^m$ Exchange Inf SupSup[Inf[\gamma > - <\phi, P_1 \gamma - u > - <\psi, P_2 \gamma - v >] \end{bmatrix}



Part. 2 Optimal Transport – Duality(DMK):
Min \gamma>
S.t.
$$P_1 \gamma = u$$

 $P_2 \gamma = v$
 $\gamma \ge 0$

$$Sup[<\phi,u> + <\psi, v>]$$
(DDMK)
$$\varphi \in IR^{m}$$

$$\psi \in IR^{n}$$

$$P_{1}^{t} \varphi + P_{2}^{t} \psi \leq C$$



Part. 2 Optimal Transport – Duality(DMK):
Min \gamma >
$$\begin{bmatrix} P_1 \gamma = U \\ P_2 \gamma = v \\ \gamma \ge 0 \end{bmatrix}$$
 $\sup [\inf [< \gamma, c - P_1^{t} \phi - P_2^{t} \psi > + < \phi, u > + < \psi, v >]]$
 $\psi \in \mathbb{R}^n$ Interpret $\sup [< \phi, u > + < \psi, v >]$ (DDMK)
 $\phi_i + \psi_j \le c_{ij} \forall (i,j)$

Kantorovich's problem:

Find a measure
$$\gamma$$
 defined on X x Y
such that $\int_{X \text{ in } X} d\gamma(x,y) = d\mu(x)$
and $\int_{Y \text{ in } Y} d\gamma(x,y) = dv(x)$

that minimizes $\iint_{X \times Y} ||x - y||^2 d\gamma(x,y)$

Dual formulation of Kantorovich's problem (Continuous):

Find two functions φ in L¹(μ) and ψ in L¹(v) Such that for all x,y, $\varphi(x) + \psi(y) \le \frac{1}{2} ||x - y||^2$ that maximize $\int_X \varphi d\mu + \int_Y \psi dv$



Kantorovich's problem:

Find a measure γ defined on X x Y such that $\int_{X \text{ in } X} d\gamma(x,y) = d\mu(x)$ and $\int_{Y \text{ in } Y} d\gamma(x,y) = dv(x)$

that minimizes $\iint_{X \times Y} || x - y ||^2 d\gamma(x,y)$

Your point of view: Try to minimize transport cost

Dual formulation of Kantorovich's problem:

Find two functions φ in L¹(μ) and ψ in L¹(v) Such that for all x,y, $\varphi(x) + \psi(y) \leq \frac{1}{2} ||x - y||^2$ that maximize $\int_X \varphi d\mu + \int_Y \psi dv$



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that minimizes $\iint_{X \times Y} || x - y ||^2 d\gamma(x,y)$

Your point of view: Try to minimize transport cost

Dual formulation of Kantorovich's problem:

Find two functions φ in $L^1(\mu)$ and ψ in $L^1(v)$ Such that for all x,y, $\varphi(x) + \psi(y) \leq \frac{1}{2} ||x - y||^2$ that maximize $\int_X \varphi d\mu + \int_Y \psi dv$

Point of view of a "transport company": Try to maximize transport price

Kantorovich's problem:

Find a measure γ defined on X x Y such that $\int_{X \text{ in } X} d\gamma(x,y) = d\mu(x)$ and $\int_{Y \text{ in } Y} d\gamma(x,y) = dv(x)$

that minimizes $\iint_{X \times Y} || x - y ||^2 d_{Y(x,y)}$

Your point of view: Try to minimize transport cost

Dual formulation of Kantorovich's problem:

Find two functions φ in L¹(μ) and ψ in L¹(v) Such that for all x,y, $\varphi(x) + \psi(y) \le \frac{1}{2} ||x - y||^2$ that maximize $\int_X \varphi(x) d\mu + \int_Y \psi(y) dv$

What they charge for loading at x



Kantorovich's problem:

Find a measure γ defined on X x Y such that $\int_{X \text{ in } X} d\gamma(x,y) = d\mu(x)$ and $\int_{Y \text{ in } Y} d\gamma(x,y) = dv(x)$

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Your point of view: Try to minimize transport cost

Dual formulation of Kantorovich's problem:

Find two functions φ in L¹(μ) and ψ in L¹(v) Such that for all x,y, $\varphi(x) + \psi(y) \le \frac{1}{2} ||x - y||^2$ that maximize $\int_X \varphi(x) d\mu + \int_Y \psi(y) dv$ What they charge for loading at x What they charge for unloading at y



Kantorovich's problem:

Find a measure γ defined on X x Y such that $\int_{X \text{ in } X} d\gamma(x,y) = d\mu(x)$ and $\int_{Y \text{ in } Y} d\gamma(x,y) = d\nu(x)$

that minimizes $\iint_{X \times Y} ||x - y||^2 d\gamma(x,y)$

Dual formulation of Kantorovich's problem:

Find two functions φ in L¹(μ) and ψ in L¹(v)

Such that for all x,y, $\varphi(x) + \psi(y) \leq \frac{1}{2} ||x - y||^2$

that maximize $\int_X \phi(x) d\mu + \int_Y \psi(y) d\nu$

Your point of view: Try to minimize transport cost

Price (loading + unloading) cannot be greater than transport cost (else you do the job yourself)

What they charge for loading at x

What they charge for unloading at y

Dual formulation of Kantorovich's problem:

Find two functions φ in L¹(μ) and ψ in L¹(v) Such that for all x,y, $\varphi(x) + \psi(y) \le \frac{1}{2} ||x - y||^2$ that maximize $\int_X \varphi(x) d\mu + \int_Y \psi(y) dv$



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If we got two functions φ and ψ that satisfy the constraint

Then it is possible to obtain a better solution by replacing ψ with φ^c defined by: For all y, $\varphi^c(y) = \inf_{x \text{ in } X} \frac{1}{2} ||x - y||^2 - \varphi(y)$



Dual formulation of Kantorovich's problem:

Find two functions φ in L¹(μ) and ψ in L¹(v) Such that for all x,y, $\varphi(x) + \psi(y) \le \frac{1}{2} ||x - y||^2$ that maximize $\int_X \varphi(x) d\mu + \int_Y \psi(y) dv$

If we got two functions ϕ and ψ that satisfy the constraint

Then it is possible to obtain a better solution by replacing ψ with φ^c defined by: For all y, $\varphi^c(y) = \inf_{x \text{ in } X} \frac{1}{2} ||x - y||^2 - \varphi(y)$

- ϕ^c is called the **c-conjugate** function of ϕ
- If there is a function φ such that $\psi = \varphi^c$ then ψ is said to be **c-concave**
- If ψ is c-concave, then $\psi^{cc} = \psi$

Dual formulation of Kantorovich's problem:

Find a c-concave function ψ

that maximizes $\int_X \psi(x) d\mu + \int_Y \psi^c(y) dv$

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Dual formulation of Kantorovich's problem:

Find a c-concave function ψ

that maximizes
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 ψ is called a **"Kantorovich potential"**

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Dual formulation of Kantorovich's problem:

Find a c-concave function ψ

that maximizes
$$\int_X \psi(x) d\mu + \int_Y \psi^c(y) dv$$

 ψ is called a **"Kantorovich potential"**

What about our initial problem ?



Dual formulation of Kantorovich's problem:

Find a c-concave function ψ

that maximizes
$$\int_X \psi(x) d\mu + \int_Y \psi^c(y) dv$$

 ψ is called a **"Kantorovich potential"**

What about our initial problem ? (i.e., this is T() that we want to find ...)



Theorem 1.

 $\forall (x,y) \in \partial_c \psi, \nabla \psi(x) - \nabla_x c(x,y) = 0$

where $\partial_c \psi = \{(x,y) | \phi(x) + \psi(y) = c(x,y)\}$ denotes the so-called c-subdifferential of ψ .



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Proof: see OTON, chap. 10.

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Proof: see OTON, chap. 10.

Heuristic argument (at the beginning of the same chapter):

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Theorem 1.

$$\forall (x,y) \in \partial_c \psi, \nabla \psi(x) - \nabla_x c(x,y) = 0$$

where $\partial_c \psi = \{(x,y) | \phi(x) + \psi(y) = c(x,y)\}$ denotes the so-called c-subdifferential of ψ .

Proof: see OTON, chap. 10.

Heuristic argument (at the beginning of the same chapter):

Consider a point (x, y) on the c-subdifferential $\partial_c \psi$, that satisfies $\phi(y) + \psi(x) = c(x, y)$ (1).





Theorem 1.

$$\forall (x,y) \in \partial_c \psi, \nabla \psi(x) - \nabla_x c(x,y) = 0$$

where $\partial_c \psi = \{(x,y) | \phi(x) + \psi(y) = c(x,y)\}$ denotes the so-called c-subdifferential of ψ .

Proof: see OTON, chap. 10.

Heuristic argument (at the beginning of the same chapter):

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The same derivation can be done with -w instead of w, and one gets:

 $\forall w, \nabla \psi(x) \cdot w = \nabla_x c(x, y) \cdot w, \text{ thus } \forall (x, y) \in \partial_c \psi, \nabla \psi(x) - \nabla_x c(x, y) = 0.$





Dual formulation of Kantorovich's problem:

Find a c-concave function ψ

that maximizes
$$\int_X \psi(x) d\mu + \int_Y \psi^c(y) dv$$

In the L_2 case, i.e. $c(x,y) = 1/2||x-y||^2$, we have $\forall (x,y) \in \partial_c \psi, \nabla \psi(x) + y - x = 0$, thus, whenever the optimal transport map T exists, we have $T(x) = x - \nabla \psi(x) = \nabla (||x||^2/2 - \psi(x))$.



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grad $\overline{\psi}(x)$ with $\overline{\psi}(x) := (\frac{1}{2} x^2 - \psi(x))$



Part. 2 Optimal Transport – convexity

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Proof.

$$\begin{split} \psi(x) &= \inf_{y} \frac{|x-y|^2}{2} - \phi(y) \\ &= \inf_{y} \frac{||x||^2}{2} - x \cdot y + \frac{||y||^2}{2} - \phi(y) \\ -\bar{\psi}(x) &= \phi(x) - \frac{||x||^2}{2} = \inf_{y} -x \cdot y + \left(\frac{||y||^2}{2} - \phi(y)\right) \\ \bar{\psi}(x) &= \sup_{y} x \cdot y - \left(\frac{||y||^2}{2} - \phi(y)\right) \end{split}$$



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Part. 2 Optimal Transport – no collision

Dual formulation of Kantorovich's problem:

Find a c-concave function ψ

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Part. 2 Optimal Transport – no collision

Dual formulation of Kantorovich's problem:

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Proof. By contradiction, suppose that you have $t \in (0, 1)$ and $x_1 \neq x_2$ such that:

$$(1-t)x_1 + tT(x_1) = (1-t)x_2 + tT(x_2)$$



Part. 2 Optimal Transport – no collision

Dual formulation of Kantorovich's problem:

Find a c-concave function ψ

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that maximizes
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Two transported particles cannot "collide"

Proof. By contradiction, suppose that you have $t \in (0, 1)$ and $x_1 \neq x_2$ such that:

$$(1-t)x_1 + tT(x_1) = (1-t)x_2 + tT(x_2)$$

$$\begin{aligned} (1-t)x_1 + t\nabla\bar{\psi}(x_1) &= (1-t)x_2 + t\nabla\bar{\psi}(x_2) \\ (1-t)(x_1 - x_2) + t(\nabla\bar{\psi}(x_1) - \nabla\bar{\psi}(x_2)) &= 0 \\ \forall v, (1-t)v \cdot (x_1 - x_2) + tv \cdot (\nabla\bar{\psi}(x_1) - \nabla\bar{\psi}(x_2)) &= 0 \\ \text{take } v &= (x_1 - x_2) \\ (1-t)\|x_1 - x_2\|^2 + t(x_1 - x_2) \cdot (\nabla\bar{\psi}(x_1) - \nabla\bar{\psi}(x_2)) &= 0 \end{aligned}$$

Part. 2 Optimal Transport – Monge-Ampere

Dual formulation of Kantorovich's problem:

Find a c-concave function ψ

that maximizes
$$\int_X \psi(x) d\mu + \int_Y \psi^c(y) dv$$

What about our initial problem ? If T(.) exists, then one can show that: T(x) = x - grad $\psi(x)$ = grad (¹/₂ x²- $\psi(x)$)

grad $\overline{\psi}(x)$ with $\overline{\psi}(x) := (\frac{1}{2} x^2 - \psi(x))$

for all borel set A, $\int_A d\mu = \int_{T(A)} (|JT|) dv$ (change of variable)

Jacobian of T (1st order derivatives)



Part. 2 Optimal Transport – Monge-Ampere

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for all borel set A,
$$\int_A d\mu = \int_{T(A)} (|JT|) dv = \int_{T(A)} (H \overline{\psi}) dv$$

Det. of the Hessian of $\overline{\psi}$ (2nd order derivatives)



Part. 2 Optimal Transport – Monge-Ampere

Dual formulation of Kantorovich's problem:

Find a c-concave function ψ

that maximizes
$$\int_X \psi(x) d\mu + \int_Y \psi^c(y) dv$$

What about our initial problem ? $T(x) = x - \text{grad } \psi(x) = \text{grad } (\frac{1}{2} x^{2} - \psi(x))$ $\text{grad } \overline{\psi}(x) \text{ with } \overline{\psi}(x) := (\frac{1}{2} x^{2} - \psi(x))$ for all borel set A, $\int_{A} d\mu = \int_{T(A)} \left(|JT| \right) dv = \int_{T(A)} \left(H \overline{\psi} \right) dv$ When μ and ν have a density u and v, $(H \overline{\psi}(x)) \cdot v(\text{grad } \overline{\psi}(x)) = u(x)$ Monge-Ampère equation

Part. 2 Optimal Transport – summary

Find a transport map T that minimizes $C(T) = \int_X ||x - T(x)||^2 d\mu(x)$



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After several rewrites and under some conditions.... (Kantorovich formulation, dual, c-convex functions)

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Solve $(H \overline{\psi}(x))$. $v(\text{grad } \overline{\psi}(x)) = u(x)$ Monge-Ampère equation (When μ and ν have a density u and v resp.)

Brenier, Mc Cann, Trudinger: *The optimal transport map is then given by:* $T(x) = \text{ grad } \overline{\psi}(x)$

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Semi-Discrete Optimal Transport



Part. 3 Optimal Transport – how to program ? (X;µ) (Y;v)



Continuous



Part. 3 Optimal Transport – how to program ? (X;µ) (Y;v)

Continuous

Semi-discrete





Part. 3 Optimal Transport – how to program ? (X;µ) (Y;v)

Continuous

Semi-discrete



Discrete







Part. 3 Optimal Transport – semi-discrete (X;µ) (Y;v)



(DMK)
$$\sup_{\psi \in \psi^c} \int_X \psi^c(x) d\mu + \int_Y \psi(y) dv$$



Part. 3 Optimal Transport – semi-discrete (X;µ) (Y;v)





$$(\text{DMK}) \quad \begin{array}{l} \underset{\psi \in \psi^c}{\overset{\text{Sup}}{\overset{\text{Sup}}{\overset{\text{Sup}}{\overset{\text{Sup}}{\overset{\text{C}}{\overset{\text{W}}{\overset{\text{C}}{\overset{\text{C}}{\overset{\text{Sup}}}{\overset{\text{Sup}}{\overset{\text{Sup}}}{\overset{\text{Sup}}{\overset{\text{Sup}}{\overset{\text{Sup}}}{\overset{\text{Sup}}{\overset{\text{Sup}}}}{\overset{\text{Sup}}}{\overset{\text{Sup}}}}}}}}}}}}}}}}}}}}} } \\ \\ \underset{{\tilde{Sup}}{\overset{\text{Sup}}}{\overset{Sup}}}}}}{\overset{Sup}}{\overset{Sup}}{\overset{Sup}}}}}}} \\ \overset{Sup}}{\overset{Sup}}{\overset{Sup}}}}} \\} \\\\} \overset{Sup}}{\overset{Sup}}}{\overset{Sup}}}}} \\\\} \overset{Sup}}{\overset{Sup}}}} \overset{Sup}}{\overset{Sup}}}}} \\\\} \overset{Sup}}{\overset{Sup}}}} \overset{Sup}}{\overset{Sup}}}} \\\\} \\\\} \overset{Sup}}}{\overset{Sup}}}} \overset{Sup}}} \overset{Sup}}} \overset{Sup}}} \overset{Sup}}{\overset{Sup}}} \overset{Sup}}} \overset{Sup}} \overset{Sup}} \overset{Sup}}} \overset{Sup}}} \overset{Sup}}} \overset{Sup}}} \overset{Sup}$$





$$\begin{array}{ll} \text{(DMK)} & \underset{\psi \in \psi^c}{\overset{\text{Sup}}{\overset{\text{Sup}}{\overset{\text{V}^c}{\overset{\text{W}^c}{(x)d\mu}}}} \int_X \psi^c(x)d\mu + \int_Y \psi(y)d\nu \\ \\ \int_X \inf_{y_j \in Y} \left[\|x - y_j\|^2 - \psi(y_j) \right] d\mu & \sum_j \psi(y_j) v_j \end{array}$$







(DMK) Sup
$$\psi \in \psi^{c}$$
 $G(\psi) = \sum_{j} \int_{\text{Lag } \psi(yj)} ||x - y_{j}||^{2} - \psi(y_{j}) d\mu + \sum_{j} \psi(y_{j}) v_{j}$

Where: Lag $\psi(yj) = \left\{ \begin{array}{cc} x & | & || & x - y_j \ ||^2 - \psi(y_j) & < || & x - y_j \ ||^2 - \psi(y_{j'}) \end{array} \right\}$ for all j' $\neq j$



Where: Lag
$$\psi(yj) = \{ x \mid || x - y_j ||^2 - \psi(y_j) < || x - y_j ||^2 - \psi(y_{j'}) \}$$
 for all $j' \neq j$

Laguerre diagram of the y_j 's (with the L₂ cost || x - y ||² used here, Power diagram)



(DMK)
$$\sup_{\psi \in \psi^{c}} G(\psi) = \sum_{j} \int_{\text{Lag } \psi(yj)} ||x - y_{j}||^{2} - \psi(y_{j}) d\mu + \sum_{j} \psi(y_{j}) v_{j}$$

Where: $\text{Lag } \psi(yj) = \left\{ x \mid ||x - y_{j}||^{2} - \psi(y_{j}) < ||x - y_{j}||^{2} - \psi(y_{j'}) \right\}$ for all $j' \neq j$
Laguerre diagram of the y_{j} 's
(with the L₂ cost $||x - y||^{2}$ used here, Power diagram)



(DMK)
$$\begin{split} & \underset{\psi \in \psi^{c}}{\operatorname{Sup}} \quad G(\psi) = \sum_{j} \int_{\text{Lag } \psi(yj)} \|x - y_{j}\|^{2} - \psi(y_{j}) \, d\mu + \sum_{j} \psi(y_{j}) \, v_{j} \\ & \text{Where: Lag } \psi(yj) = \left\{ \begin{array}{c} x \mid \|x - y_{j}\|^{2} - \psi(y_{j}) < \|x - y_{j}\|^{2} - \psi(y_{j'}) \end{array} \right\} \text{ for all } j' \neq j \\ & & & \\ & & \text{Laguerre diagram of the } y_{j}'s \\ & & (\text{with the } L_{2} \operatorname{cost} \||x - y\|^{2} \operatorname{used here, Power diagram}) \end{split}$$



 ψ is determined by the weight vector $[\psi(y_1)\,\psi(y_2)\,\ldots\,\psi(y_m)]$



Part. 3 Power Diagrams

Voronoi diagram: $Vor(x_i) = \{ x \mid d^2(x, x_i) < d^2(x, x_i) \}$

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Part. 3 Power Diagrams

Voronoi diagram: $Vor(x_i) = \{ x \mid d^2(x, x_i) < d^2(x, x_j) \}$

Power diagram: $Pow(x_i) = \{ x \mid d^2(x, x_i) - \psi_i < d^2(x, x_j) - \psi_j \}$

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Part. 3 Power Diagrams



Part. 3 Optimal Transport



<u>Theorem</u>: (direct consequence of MK duality alternative proof in [Aurenhammer, Hoffmann, Aronov 98]):

Given a measure μ with density, a set of points (y_j) , a set of positive coefficients v_j such that $\sum v_j = \int d\mu(x)$, it is possible to find the weights $W = [\psi(y_1) \ \psi(y_2) \ \dots \ \psi(y_m)]$ such that the map $T_S{}^W$ is the unique optimal transport map between μ and $v = \sum v_j \ \delta(y_j)$


Part. 3 Optimal Transport



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Proof:
$$G(\psi) = \sum_{j} \int_{\text{Lag } \psi(yj)} ||x - y_j||^2 - \psi(y_j) d\mu + \sum_{j} \psi(y_j) v_j$$

Is a concave function of the weight vector $[\psi(y_1) \psi(y_2) \dots \psi(y_m)]$



Part. 3 Optimal Transport



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Idea of the proof

Consider the function
$$f_T(W) = \int (||x - T(x)||^2 - \psi(T(X))) d\mu(x)$$

The (unknown) weights $W = [\Psi(y_1) \Psi(y_2) \dots \Psi(y_m)]$





Idea of the proof

Consider the function

$$f_{T}(W) = \int (||x - T(x)||^{2} - \psi(T(X))) d\mu(x)$$



T : an arbitrary but fixed assignment.









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$$f_T(W) = \int (||x - T(x)||^2 - \psi(T(X))) d\mu(x)$$









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 $f_T(W)$ is linear in W f: W $\rightarrow f_{T_W}(W)$ is **CONCAVE !!** (because its graph is the lower enveloppe of linear functions)





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Consider
$$g(W) = f_{T_W}(W) + \sum v_j \psi_j$$



Idea of the proof

Consider the function
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$$f_{T}(W)$$

 $f_{T}(W)$
 $f_{T_{W}}(W) = \min_{T} f_{T}(W)$
 W

Consider
$$g(W) = f_{T_W}(W) + \sum v_j \psi_j$$

 $\partial g / \partial \Psi_{j} = V_{j} - \int_{Lag^{\Psi}(yj)} d\mu(x)$ and g is concave.

Semi-discrete OT Summary:

(DMK) $\begin{array}{c} \text{Sup} \\ \psi \in \psi^{c} \end{array} \quad G(\psi) = \int_{X} \psi^{c}(x) d\mu + \int_{Y} \psi(y) dv \end{array}$

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$$G(\psi) = g(W) = \sum_{j} \int_{Lag^{\psi}(yj)} ||x - y_j||^2 - \psi(y_j) d\mu + \sum_{j} \psi(y_j) v_j \text{ is concave}$$



Semi-discrete OT Summary:

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$$\partial G / \partial \Psi_j = V_j - \int_{Lag(yj)} d\mu(x)$$
 (= 0 at the maximum)



Semi-discrete OT Summary:

(DMK) $\begin{array}{c} \text{Sup} \\ \psi \in \psi^c \end{array} \quad G(\psi) = \int_X \psi^c(x) d\mu + \int_Y \psi(y) dv \end{array}$

$$G(\psi) = g(W) = \sum_{j} \int_{Lag^{\psi}(yj)} ||x - y_j||^2 - \psi(y_j) d\mu + \sum_{j} \psi(y_j) v_j \text{ is concave}$$



Desired mass at y_i

Mass transported to y_i



$$\partial G / \partial \Psi_{j} = V_{j} - \int_{Lag(yj)} d\mu(x)$$

$$\partial G / \partial \Psi_{j} = V_{j} - \int_{Lag(yj)} d\mu(x)$$

$$\partial^{2}G / \partial \Psi_{i}\Psi_{j} = - \partial / \partial \Psi_{j} \int_{Lag(yj)} d\mu(x)$$

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$$\partial G / \partial \Psi_{j} = V_{j} - \int_{\text{Lag}(yj)} d\mu(x)$$
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$$\Psi_{j} \leftarrow \Psi_{j+} \delta \Psi_{j}$$

уj

$$\partial G / \partial \Psi_{j} = V_{j} - \int_{Lag(yj)} d\mu(x)$$

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$$\Psi_{j} \leftarrow \Psi_{j+} \delta \Psi_{j}$$

Reynold's thm:





Reynold's thm:



Reynold's thm:



Reynold's thm:

$$\partial_{\partial_i} \psi_j \int_{\text{Lag}(yj)} d\mu(x) = \int_{\partial_i \text{Lag}(yj)} v.n \ d\mu(x)$$

 $f_{ij}(x) = 0$

$$\begin{aligned} \mathsf{c}(\mathsf{x},\mathsf{y}_{\mathsf{i}}) &- \mathsf{c}(\mathsf{x},\mathsf{y}_{\mathsf{j}}) + \psi_{\mathsf{j}} - \psi_{\mathsf{i}} &= 0 \\ \mathsf{df}_{\mathsf{ij}} &= \mathsf{grad}_{\mathsf{x}}(\mathsf{c}(\mathsf{x},\mathsf{y}_{\mathsf{i}}) - \mathsf{c}(\mathsf{x},\mathsf{y}_{\mathsf{j}}))\mathsf{d}\mathsf{x} + \mathsf{d}\psi_{\mathsf{j}} \end{aligned}$$





Reynold's thm:

$$\partial_{\partial_i} \psi_j \int_{\text{Lag}(yj)} d\mu(x) = \int_{\partial_i \text{Lag}(yj)} v.n \ d\mu(x)$$

 $f_{ij}(x) = 0$

$$\begin{split} & \mathsf{c}(\mathsf{x},\mathsf{y}_i) - \mathsf{c}(\mathsf{x},\mathsf{y}_j) + \psi_j - \psi_i = 0 \\ & \mathsf{d} \mathsf{f}_{ij} = \mathsf{grad}_\mathsf{x}(\mathsf{c}(\mathsf{x},\mathsf{y}_i) - \mathsf{c}(\mathsf{x},\mathsf{y}_j))\mathsf{d} \mathsf{x} + \mathsf{d} \psi_j \\ & \delta \mathsf{x} = \delta \mathsf{h} \ \mathsf{n} = \delta \mathsf{h} \ \mathsf{grad}_\mathsf{x} \ \mathsf{f}_{ij}(\mathsf{x}) \, / \, || \ \mathsf{grad}_\mathsf{x} \ \mathsf{f}_{ij}(\mathsf{x}) || \end{split}$$





Reynold's thm:

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Reynold's thm:

$$\partial_{\partial_i} \psi_j \int_{\text{Lag}(yj)} d\mu(x) = \int_{\partial_i \text{Lag}(yj)} v.n \ d\mu(x)$$

 $f_{ij}(x) = 0$

$$c(x,y_{i}) - c(x,y_{j}) + \psi_{j} - \psi_{i} = 0$$

$$df_{ij} = \operatorname{grad}_{x}(c(x,y_{i}) - c(x,y_{j}))dx + d\psi_{j}$$

$$\delta x = \delta h \ n = \delta h \ \operatorname{grad}_{x} f_{ij}(x) / || \ \operatorname{grad}_{x} f_{ij}(x) ||$$

$$\partial h / \partial \psi_j = -1/ \parallel \operatorname{grad}_x c(x, y_i) - \operatorname{grad}_x c(x, y_j) \parallel$$





Reynold's thm:

$$\partial_{\partial_i} \psi_j \int_{\text{Lag}(yj)} d\mu(x) = \int_{\partial_i \text{Lag}(yj)} v.n \ d\mu(x)$$

 $f_{ij}(x) = 0$

$$c(x,y_{i}) - c(x,y_{j}) + \psi_{j} - \psi_{i} = 0$$

$$df_{ij} = \operatorname{grad}_{x}(c(x,y_{i}) - c(x,y_{j}))dx + d\psi_{j}$$

$$\delta x = \delta h n = \delta h \operatorname{grad}_{x} f_{ij}(x) / || \operatorname{grad}_{x} f_{ij}(x)||$$

$$\partial h / \partial \psi_{j} = -1 / || \operatorname{grad}_{x} c(x,y_{i}) - \operatorname{grad}_{x} c(x,y_{j}) ||$$

$$\partial_{\lambda} \psi_{j} \int_{\text{Lag}(yj)} d\mu(x) = \int_{\text{Lag}(yi) \cap \text{Lag}(yj)} -1/|| \operatorname{grad}_{x} c(x,y_{i}) - \operatorname{grad}_{x} c(x,y_{j}) || d\mu(x)$$



$$\partial^2 / \partial \psi_i \partial \psi_j F = \int_{\text{Lag}(yi) \cap \text{Lag}(yj)} -1/|| \text{ grad}_x c(x, y_i) - \text{grad}_x c(x, y_j) || d\mu(x)$$

$$\partial^2$$
 / ∂^2 / $\partial^2 \psi_i^2 F$ = - $\sum \partial^2$ / ∂^2 / $\partial^2 \psi_i \partial^2 \psi_j$



$$\partial^2 / \partial \psi_i \partial \psi_j F = \int_{\text{Lag}(y_i) \cap \text{Lag}(y_j)} -1/|| \operatorname{grad}_x c(x, y_i) - \operatorname{grad}_x c(x, y_j) || d\mu(x)$$

$$\partial^2$$
 / ∂^2 / ∂^ψ i 2F = - \sum ∂^2 / ∂^ψ i ∂^ψ j

$$\begin{split} c(x,y) &= || \ x - y \ ||^2 \\ \partial^2 / \partial \psi_{i \,\partial} \psi_j F = \int_{\text{Lag}(yi) \,\cap \, \text{Lag}(yj)} 1 \ / \, || \ xj - xi \, || \ d\mu(x) \end{split}$$

$$\partial^2 / \partial \psi_i \partial \psi_j F = \int_{\text{Lag}(y_i) \cap \text{Lag}(y_j)} -1/|| \operatorname{grad}_x c(x, y_i) - \operatorname{grad}_x c(x, y_j) || d\mu(x)$$

$$\partial^2$$
 / ∂^2 / ∂^{ψ} i 2F = - \sum ∂^2 / ∂^{ψ} i ∂^{ψ} j

$$\begin{split} c(x,y) &= || \ x - y \ ||^2 \\ \partial^2 / \partial \psi_{i \,\partial} \psi_j F &= \int_{\text{Lag}(yi) \,\cap \, \text{Lag}(yj)} 1 \ / \, || \ xj - xi \ || \ d\mu(x) \\ IP_I \ \text{FEM Laplacian (not a big surprise)} \end{split}$$

Part. 3 Optimal Transport Let's program it !

Hierarchical algorithm [Mérigot] Geometry, 3D [L], [L, Schwindt] Damped Newton algorithm, [Kitagawa, Mérigot, Thibert]





Optimal Transport applications in computational physics

Innia


Hamilton, Legendre, Maupertuis

Lagrange



The Least Action Principle

Axiom 1: There exists a function L(x, x, t)

that describes the state of a physical system

Short summary of the 1st chapter of Landau, Lifshitz Course of Theoretical Physics





Hamilton, Legendre, Maupertuis

Lagrange



The Least Action Principle

Axiom 1: There exists a function L(x, x, t)

position

that describes the state of a physical system





Hamilton, Legendre, Maupertuis

Lagrange



The Least Action Principle

Axiom 1: There exists a function L(x, x, t)



that describes the state of a physical system





Hamilton, Legendre, Maupertuis

Lagrange



The Least Action Principle

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that describes the state of a physical system





Hamilton, Legendre, Maupertuis

Lagrange



The Least Action Principle

Axiom 1: There exists a function L(x, x, t)

that describes the state of a physical system

Axiom 2: The movement (time evolution) of the physical system minimizes the following integral

$$\int_{t_1}^{t_2} L(x, x, t) dt$$





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that describes the state of a physical system

Axiom 2: The movement (time evolution) of the physical system minimizes the following integral

$$\int_{t_1}^{t_2} L(\mathbf{x}, \mathbf{x}, t) dt$$

L(x,x,t) dt

Axiom 1: There exists a function L(x, x, t)

that describes the state of a physical system

=

Axiom 2: The movement (time evolution) of the physical system minimizes the following integral

Theorem 1: (Lagrange equation):

$$\frac{\partial L}{\partial x} = \frac{d}{dt} \frac{\partial L}{\partial \dot{x}}$$



Axiom 1: There exists L Axiom 2: The movement minimizes

$$\int_{t_1}^{t_2} L(x, x, t) dt$$

Axiom 3:

Invariance w.r.t. change of Gallileo frame + hom. + isotrop. :

$$\frac{x'}{t'} = \frac{x+vt}{t}$$

Theorem 1: (Lagrange equation):

$$\frac{\partial L}{\partial x} = \frac{d}{dt} \frac{\partial L}{\partial \dot{x}}$$



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Theorem 1: (Lagrange equation):

$$\frac{\partial L}{\partial x} = \frac{d}{dt} \frac{\partial L}{\partial \dot{x}}$$

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Axiom 3:

Invariance w.r.t. change of Gallileo frame + hom. + isotrop. :

$$\frac{x'}{t'} = \frac{x + vt}{t}$$

Theorem 2:

$$\mathbf{\hat{x}} \frac{\partial \mathbf{L}}{\partial \mathbf{\hat{x}}} - \mathbf{L} = \mathsf{cte}$$

Axiom 1: There exists L Axiom 2: The movement minimizes

$$\int_{t_1}^{t_2} L(x, x, t) dt$$

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Theorem 2:

$$\overset{\bullet}{\mathbf{x}} \frac{\partial \mathbf{L}}{\partial \overset{\bullet}{\mathbf{x}}} - \mathbf{L} = \mathbf{cte}$$

Homogeneity of time \rightarrow Preservation of **energy**

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Homogeneity of time \rightarrow Preservation of **energy**

Homogeneity of space \rightarrow Preservation of **momentum**



Axiom 1: There exists L Axiom 2: The movement minimizes

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Theorem 1: (Lagrange equation):

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Axiom 3:

Invariance w.r.t. change of Gallileo frame + hom. + isotrop. :

$$\frac{x'}{t'} = \frac{x+vt}{t}$$

Theorem 2:

$$\frac{\partial L}{\partial \mathbf{x}} - L = cte$$

Homogeneity of time \rightarrow Preservation of **energy**

Homogeneity of space \rightarrow Preservation of **momentum**



Axiom 3:

Invariance w.r.t. change of

Axiom 1: There exists L Axiom 2: The movement minimizes



Axiom 1: There exists L

Axiom 2: The movement minimizes $\int L$

Theorem 1: (Lagrange equation):

$$\frac{\partial L}{\partial x} = \frac{d}{dt} \frac{\partial L}{\partial x}$$

Free particle:

Theorem 3: v = cte (Newton's law I)

Axiom 3:

Invariance w.r.t. change of Gallileo frame + hom. + isotrop. :

$$\frac{x'}{t'} = \frac{x+vt}{t}$$

Theorem 2:

OX

$$\frac{1}{x}\frac{\partial L}{\partial L}$$
 - L = cte

Homogeneity of time \rightarrow Preservation of **energy**

Homogeneity of space \rightarrow Preservation of **momentum**



Axiom 1: There exists L

Axiom 2: The movement minimizes $\int L$

Theorem 1: (Lagrange equation):

$$\frac{\partial L}{\partial x} = \frac{d}{dt} \frac{\partial L}{\partial x}$$

Free particle:

Theorem 3: v = cte (Newton's law I)

Expression of the Lagrangian:

 $L = \frac{1}{2} m v^2$

Axiom 3:

Invariance w.r.t. change of Gallileo frame + hom. + isotrop. :

$$\frac{x'}{t'} = \frac{x+vt}{t}$$

Theorem 2:

$$\overset{\bullet}{\mathbf{x}} \frac{\partial \mathbf{L}}{\partial \overset{\bullet}{\mathbf{x}}} - \mathbf{L} = \mathbf{cte}$$

Homogeneity of time \rightarrow Preservation of **energy**

Homogeneity of space \rightarrow Preservation of **momentum**



Axiom 1: There exists L

Axiom 2: The movement minimizes $\int L$

Theorem 1: (Lagrange equation):

$$\frac{\partial L}{\partial x} = \frac{d}{dt} \frac{\partial L}{\partial x}$$

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Expression of the Lagrangian:

 $L = \frac{1}{2} m v^2$

Expression of the Energy:

 $E = \frac{1}{2} m v^2$

Axiom 3:

Invariance w.r.t. change of Gallileo frame + hom. + isotrop. :

$$\frac{x'}{t'} = \frac{x+vt}{t}$$

Theorem 2:

$$\frac{\partial L}{\partial \mathbf{X}} - L = cte$$

Homogeneity of time \rightarrow Preservation of **energy**

Homogeneity of space \rightarrow Preservation of **momentum**



Axiom 1: There exists L

Axiom 2: The movement minimizes $\int L$

Theorem 1: (Lagrange equation):

$$\frac{\partial L}{\partial x} = \frac{d}{dt} \frac{\partial L}{\partial x}$$

Free particle:

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Expression of the Lagrangian:

 $L = \frac{1}{2} m v^2$

Expression of the Energy:

$$E = \frac{1}{2} m v^2$$

Invariance w.r.t. change of Gallileo frame + hom. + isotrop. :

$$\frac{x'}{t'} = \frac{x+vt}{t}$$

Particle in a field:

Expression of the Lagrangian: $L = \frac{1}{2} m v^2 - U(x)$



Axiom 1: There exists L

Axiom 2: The movement minimizes $\int L$

Theorem 1: (Lagrange equation):

$$\frac{\partial L}{\partial x} = \frac{d}{dt} \frac{\partial L}{\partial x}$$

Free particle:

Theorem 3: v = cte (Newton's law I)

Expression of the Lagrangian:

 $L = \frac{1}{2} m v^2$

Expression of the Energy:

 $E = \frac{1}{2} m v^2$

Axiom 3:

Invariance w.r.t. change of Gallileo frame + hom. + isotrop. :

$$\frac{x'}{t} = \frac{x+vt}{t}$$

Particle in a field:

Expression of the Lagrangian: $L = \frac{1}{2} m v^2 - U(x)$ Expression of the Energy: $E = \frac{1}{2} m v^2 + U(x)$

Axiom 1: There exists L

Axiom 2: The movement minimizes $\int L$

Theorem 1: (Lagrange equation):

$$\frac{\partial L}{\partial x} = \frac{d}{dt} \frac{\partial L}{\partial x}$$

Free particle:

Theorem 3: v = cte (Newton's law I)

Expression of the Lagrangian:

 $L = \frac{1}{2} m v^2$ Expression of the Energy:

 $E = \frac{1}{2} m v^2$

Axiom 3:

Invariance w.r.t. change of Gallileo frame + hom. + isotrop. :

$$\frac{x}{t} = \frac{x+vt}{t}$$

Particle in a field:

Expression of the Lagrangian: $L = \frac{1}{2} m v^{2} - U(x)$ Expression of the Energy: $E = \frac{1}{2} m v^{2} + U(x)$ Theorem 4: $mx = -\nabla U \text{ (Newton's law II)}$

The Least Action Principle (relativistic setting – just for fun...)

Axiom 1: There exists L

Axiom 2: The movement minimizes $\int L$

Theorem 1: (Lagrange equation):

$$\frac{\partial L}{\partial x} = \frac{d}{dt} \frac{\partial L}{\partial \dot{x}}$$

Axiom 3:

Invariance w.r.t. Lorentz change of frame

$$\begin{aligned} x' \\ t' &= \frac{(x-vt) \times \gamma}{(t-vx/c^2) \times \gamma} \\ \gamma &= 1 / \sqrt{(1-v^2/c^2)} \end{aligned}$$



The Least Action Principle (relativistic setting – just for fun...)

Axiom 1: There exists L

Axiom 2: The movement minimizes $\int L$

Theorem 1: (Lagrange equation):

$$\frac{\partial L}{\partial x} = \frac{d}{dt} \frac{\partial L}{\partial \dot{x}}$$

Axiom 3:

Invariance w.r.t. Lorentz change of frame

$$\gamma = 1 / \sqrt{(1 - v^2 / c^2)}$$

Theorem 5:

$$E = \frac{1}{2} \gamma m v^2 + mc^2$$



(quantum physics setting – just for fun...)

In quantum mechanics non just the extreme path contributes to the probability amplitude

$$K(B, A) = \sum_{over all possible paths} \phi[x(t)]$$

where $\Phi[x(t)] = A \exp\left\{\frac{i}{\hbar}S[x(t)]\right\}$

Feynman's path integral formula

$$K(B,A) = \int_{A}^{B} \exp\left(\frac{i}{\hbar}S[B,A]Dx(t)\right)$$

 $P(B, A) = |K(2, 1)|^2$







Fluids – Benamou Brenier







 ρ_2



Fluids – Benamou Brenier



Minimize

$$A(\rho,v) = (t_2-t_1) \int_{t_1}^{t_2} \int_{\Omega} \rho(x,t) ||v(t,x)||^2 dx dt$$

$$s.t. \ \rho(t_1,.) = \rho_1 \quad ; \quad \rho(t_2,.) = \rho_2 \quad ; \quad \frac{d \ \rho}{dt} = - \ div(\rho v)$$



Fluids – Benamou Brenier



$$\begin{array}{l} \text{Minimize} \\ A(\rho, v) = (t_2 - t_1) \\ t_1 \\ \Omega \\ \text{s.t. } \rho(t_1, .) = \rho_1 \\ \text{s.t. } \rho(t_2, .) = \rho_2 \\ \text{;} \\ \begin{array}{l} \frac{d \ \rho}{dt} = - \ \text{div}(\rho v) \end{array} \end{array} \\ \begin{array}{l} \text{Minimize } C(T) = \\ \int_{\Omega} \rho_1(x) \left\| x - T(x) \right\| 2 \ dx \\ \text{s.t. } T \text{ is measure-preserving} \end{array}$$





Le schéma [Mérigot-Gallouet] Applications en graphisme: [De Goes et.al] (power particles)



Part. 4 Optimal Transport – Fluids Let's code it !





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Vortices in "ether" ?

René Descartes - 1663







COBE 1992

The Data: #1: the Cosmic Microwave Background



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The Data: #1 the Cosmic Microwave Background



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The Data: #2 redshift acquisition surveys





The Data: #2 redshift acquisition surveys




Part. 4 To infinity and beyond...

The Data: #2 redshift acquisition surveys





Part. 4 To infinity and beyond... pc/h : parsec (= 3.2 années lumières)



The millenium simulation project, Max Planck Institute fur Astrophysik



Part. 4 To infinity and beyond...

The universal swimming pool





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Time = Now





Time = BigBang

(-13.7 billion Y)



"Time-warped" map of the universe





Cosmic Microware Background: "Fossil light" emitted 380 000 Y after BigBang and measured now

Coop. with MOKAPLAN & Institut d'Astrophysique de Paris & Observatoire de Paris

"Time-warped" map of the universe



Cosmic Microware Background: "Fossil light" emitted 380 000 Y after BigBang and measured now

Do they match ?

"Time-warped" map of the universe



Conclusions Open Questions References Online resources

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Conclusions – Open questions

* Connections with physics, Legendre transform and entropy ?

[Cuturi & Peyré] – regularized discrete optimal transport – why does it work ? Hint 1: Minimum action principle subject to conservation laws Hint 2: Entropy = dual of temperature ; Legendre = Fourier[(+,*) → (Max,+)]...

* More continuous numerical algorithms ? [Benamou & Brenier] fluid dynamics point of view – very elegant, but 4D problem !! FEM-type adaptive discretization of the subdifferential (graph of T) ?

* Can we characterize OT in other semi-discrete settings ? measures supported on unions of spheres piecewise linear densities

* Connections with computational geometry ?

Singularity set **[Figalli]** = set of points where T is discontinuous Looks like a "mutual power diagram", anisotropic Voronoi diagrams

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Conclusions - References

A Multiscale Approach to Optimal Transport, **Quentin Mérigot**, Computer Graphics Forum, 2011

Variational Principles for Minkowski Type Problems, Discrete Optimal Transport, and Discrete Monge-Ampere Equations Xianfeng Gu, Feng Luo, Jian Sun, S.-T. Yau, ArXiv 2013

Minkowski-type theorems and least-squares clustering **AHA! (Aurenhammer, Hoffmann, and Aronov),** SIAM J. on math. ana. 1998

Topics on Optimal Transportation, 2003 Optimal Transport Old and New, 2008 Cédric Villani



Conclusions - References

Polar factorization and monotone rearrangement of vector-valued functions **Yann Brenier**, Comm. On Pure and Applied Mathematics, June 1991

A computational fluid mechanics solution of the Monge-Kantorovich mass transfer problem, **J.-D. Benamou**, **Y. Brenier**, Numer. Math. 84 (2000), pp. 375-393

Pogorelov, Alexandrov – Gradient maps, Minkovsky problem (older than AHA paper, some overlap, in slightly different context, formalism used by Gu & Yau)

Rockafeller – Convex optimization – Theorem to switch inf(sup()) – sup(inf()) with convex functions (used to justify Kantorovich duality)

Filippo Santambrogio – Optimal Transport for Applied Mathematician, Calculus of Variations, PDEs and Modeling – Jan 15, 2015

Gabriel Peyré, Marco Cuturi, Computational Optimal Transport, 2018



Online resources

All the sourcecode/documentation available from: http://alice.loria.fr/software/geogram Demo: www.loria.fr/~levy/GLUP/vorpaview

 * L., A numerical algorithm for semi-discrete L2 OT in 3D, ESAIM Math. Modeling and Analysis, 2015

* L. and E. Schwindt, Notions of OT and how to implement them on a computer, Computer and Graphics, 2018.



J. Lévy – 1936-2018 - To you Dad, I miss you so much.

Bonus Slides The Isoperimetric Inequality

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For a given volume, ball is the shape that minimizes border area



L₁ **Sobolev inegality:** Given f: $IR^n \rightarrow IR$ sufficiently regular

$$\int |\operatorname{grad} f| \ge n \operatorname{Vol}(B_2^n)^{1/n} \left(\int f^{n/(n-1)} \right) (n-1)/n$$

Explanation in [Dario Cordero Erauquin] course notes



L₁ **Sobolev inegality:** Given f: $IR^n \rightarrow IR$ sufficiently regular

$\int |\operatorname{grad} f| \ge n \operatorname{Vol}(B_2^n)^{1/n} \left(\int f^{n/(n-1)} \right)^{(n-1)/n}$

Explanation in [Dario Cordero Erauquin] course notes



L₁ **Sobolev inegality:** Given f: $IR^n \rightarrow IR$ sufficiently regular

$$\int |\operatorname{grad} f| \ge n \operatorname{Vol}(B_2^n)^{1/n} \left(\int f^{n/(n-1)} \right) (n-1)/n$$

Explanation in [Dario Cordero Erauquin] course notes



L₁ **Sobolev inegality:** Given f: $IR^n \rightarrow IR$ sufficiently regular

Consider a compact set Ω such that Vol(Ω) = Vol(B₂³) and f = the indicatrix function of Ω

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L1 Sobolev inegality: a proof with OT [Gromov]

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 $| \text{grad } f | \ge n \text{ Vol}(B_2^n)^{1/n}$



Bonus Slides Plotting the potential & optics

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The [AHA] paper summary:

- The optimal weights minimize a convex function
- The gradient and Hessian of this convex function is easy to compute

Note: the weight w(s) correspond to the Kantorovich potential $\psi(x)$ (solves a "discrete Monge-Ampere" equation)

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converging beams" con compensate the cos(x) expansion by "re-concentrating" the paints



Plotting the potential, "optics" $d^{2}(p_{i,q}) - w_{i}^{+} \leq d^{2}(p_{i,q}) - w_{j}^{+} V_{j}^{-}$ $d^{2}(p_{i}, q-T) < d^{2}(j, q-T)$ V, $(p_i - q + T)^2 \leq (p_i - q + T)^2$ $d^{2}(p_{i},q) + 2T.(p_{i}-q) + T^{2} \leq d^{2}(p_{j},q) + 2T.(p_{j}-q) + T^{2} \vee_{j}$ d²(pi,q) + 2T.pi <d²(pj,q) +2T.pj $W_i^2 = -2T \cdot p_i'$ + che hi? = (2Tipit Cite) = hi= VZ(T-pi - min(T-p)) Translation d'un diagramme de bionoi sectionnel-Delivement en racine carres-

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Plotting the potential, "optics" Numerical Experiment: A disk becomes two disks

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